A MATHEMATICAL MODEL OF BLOCH NMR FLOW EQUATION FOR FIELD CYCLING TISSUE IMAGING

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INTRODUCTION

Field-cycling magnetic resonance imaging (FC MRI) allows switching of the magnetic field during an imaging scan¹. FC MRI has been very successful in relaxometry studies and there arises the need to offer more theoretical supports to the rich experimental results available in NMR laboratories. These theories are expected to offer new ways of interpreting the results for new discoveries. In view of this, we have developed a methodology based on the time – independent Bloch NMR flow equations for calculating the transverse magnetization in terms of the applied RF field.

MATHEMATICAL FORMULATION

In this study, we shall consider fluid whose spins are in a motion in which the transverse magnetization does not change appreciably with time of motion^{2, 3}. In a rotating frame of reference, we shall assume that under the influence of RF magnetic field as derived in the earlier studies², resonance condition exists at Larmor frequency^{2,3}: $f_{\alpha} = \gamma B - \omega = 0$

From the Bloch equations, the following equation as been derived^{2,3}: $v^2 \frac{d^2 M_y}{dx^2} + v \left(\frac{1}{T_1} + \frac{1}{T_2} + \frac{dv}{dx}\right) \frac{dM_y}{dx} + \left(\gamma^2 B_1^2(x) + \frac{1}{T_1T_2}\right) M_y = \frac{M_y \mathcal{B}_1(x)}{T_1}$ (1)

B₁ is the spatially varying RF magnetic field and v is the spatial fluid flow velocity. We shall assume that the variable velocity is given as follows: $v(x) = \frac{x}{\delta}$ (2)

Therefore, if we sample the MRI signal at the point where the transverse magnetization has the largest magnitude ($M_0 \approx 0$), we have:

$$x^{2} \frac{d^{2}M_{y}}{dx^{2}} + x(1 + \delta T_{0})\frac{dM_{y}}{dx} + (\gamma^{2}B_{1}^{2}(x)\delta^{2} + \delta^{2}T_{g})M_{y} = 0; T_{0} = \frac{1}{T_{1}} + \frac{1}{T_{2}}, T_{g} = \frac{1}{T_{1}T_{2}}$$
(3)

In the presence of magnetic field gradient, we shall write the spatially dependent frequency as follows: $\omega(x) = \omega_0 + \omega_1 = \gamma B_0 + \gamma G x$ and $\alpha = \omega - \omega_0$ (4)

Using Eqn (9), we can easily convert spatial information into frequency information such that: $\alpha^{2} \frac{d^{2}M_{y}}{d\alpha^{2}} + (1 + \delta T_{0})\alpha \frac{dM_{y}}{d\alpha} + (\alpha^{2}\delta^{2} + \delta^{2}T_{g})M_{y} = 0$ (5)

Eqn (16) is an equation transformable to Bessel equation, whose solution is given as: $M_{y} = \alpha^{-\frac{\delta T_{0}}{2}} \left[C_{1}J_{n}(\delta\alpha) + C_{2}Y_{n}(\delta\alpha) \right] = (\omega_{1})^{-\frac{\delta T_{0}}{2}} \left[C_{1}J_{n}(\delta\omega_{1}) + C_{2}Y_{n}(\delta\omega_{1}) \right]$ (6)

where $n = \delta \frac{\sqrt{T_0^2 - 4T_g}}{2}$, J_n and Y_n are Bessel functions of the first and second kind respectively; C_1 and C_2 are constants. However, since the NMR signal must have finite value δT_0

even when the RF B₁ field is removed, C₂ = 0, and then we have:
$$M_y = C_1(\omega_1)^{-\frac{1}{2}} J_n(\delta\omega_1)$$
 (7)

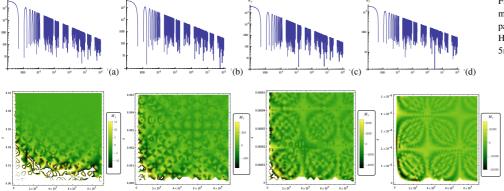


Fig. 1: 2D Plots of the NMR transverse magnetization as function of ω_1 using relaxation parameters (at 1.5T) of ⁴ (a) Skeletal muscle (b) Heart muscle (c) Liver (d) Kidney. We used $\delta = 5$ ms and $C_1 = 5 \times 10^6$.

Fig. 2: 3D images of the NMR transverse magnetization as function of ω_1 and the pulse time δ for kidney at 1.5T⁴. We have shown the behaviour of M_y at different ranges of δ and T₁ = 0.830s, T₂ = 0.082s, C₁ = 5 × 10⁶.

DISCUSSIONS AND CONCLUSION

Using Eqn (7) and the relaxation times of selected human tissues, we have shown that it is possible to do computational imaging with relaxometry data. Fig. 1 illustrates spectroscopic capabilities of the results we have obtained in this study. The tissues we considered showed unique peaks which correspond to unique values of cycling RF B₁ fields. Fig. 2 shows that the field cycling process can be easily transformed into 3D tissue mapping and it is quite interesting to see that patterns of the signals are very different as the pulse time changes. We also observed that 3D mapping shows very unique signal magnitude and slightly different patterns for different tissues at 1.5T. In conclusion, we see that we can easily use the results in this study o show contrast between various tissues and same tissues with changing T_2 values. We can also use the results to map the changes in tissue molecular dynamics at higher RF field values.

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REFERENCES

- Hógáin DO, Davies GR, Baroni S, Aime S and Lurie DJ. The use of contrast agents with fast field-cycling magnetic resonance imaging. Phys. Med. Biol. 2011; 56: 105.
 Awojoyogbe O.B., Dada O.M., Faromika O.P., Dada O.E. Mathematical Concept of the Bloch Flow Equations for General Magnetic Resonance Imaging: A Review; Concepts in Magnetic Resonance Part A, 2011; 38A (3): 85–101.
- Awojoyogbe OB, Dada M. Mathematical design of a magnetic resonance imaging sequence based on bloch NMR flow equations and Bessel functions. Chin J Magn Reson Imaging. 2013; 5(1):321-270.
- 4. Brix G, Kolem H, Nitz WR, Bock M, Huppertz A, Zech CJ, and Dietrich O. Basics of MRI and Magnetic Resonance Spectroscopy. Springer Berlin Heidelberg, 2008, pp 17.